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## DISPERSION OF THE SPEED OF SOUND IN WATER-ORGANIC SOLUTIONS

**Abstract:** Concentration and temperature studies of spontaneous Mandelstam-Brillouin scattering of light in aqueous solutions of *g*-picoline showed the manifestation of negative dispersion of hypersound velocity at concentrations  $C \sim 0.2$  mole fractions of *g*-picoline. The temperature and concentration dependence of the negative dispersion is associated with the process of solution structurization, which reaches its optimum value at a concentration of  $C = 0.06$  mole fractions of *g*-picoline.

**Key words:** dispersion, relaxation of bulk viscosity, hypersound, relaxation time, water-organic solution.

**Language:** English

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### Introduction

The work is devoted to the description of an experiment on measuring the temperature dependences of the hypersound velocity at different frequencies in water-organic solutions with singular points. Interest in the study of such solutions is due to the following. As is known (see, for example, monograph [1]), there is a class of water-organic solutions in which an anomalously high scattering peak unrelated to molecular scattering is observed in

a certain fairly narrow range of concentrations. The physical nature of this peak is currently being actively discussed (see, for example, [2] and references cited in this work). It is believed that the main reason for the appearance of this peak is the appearance of a supramolecular structure in a certain range of temperatures and concentrations. It is necessary to mention the previous works [3 - 6], in which it has been shown that stable liquid molecular complexes with sizes of 100 - 200 nm spontaneously arise in

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aqueous solutions of polar organics, the density of which differs from the density of the surrounding liquid. It has been shown in [7 - 12] that anomalous light scattering is associated with the formation of a nanodroplet phase in water-organic mixtures.

In this work, a mixture of "water - g-picoline" is studied. Interest in this particular mixture is due to the fact that in this solution in a certain concentration range there is an abnormally high light scattering peak at a content of g-picoline  $\gg 0.06$  mole fractions (ppm), see Fig. III.14 in the monograph [1] and comments on this figure. At the same time, in accordance with the literature data (see, for example, [11]), g-picoline has infinite solubility in water, i.e. the mixture "g-picoline - water" does not separate at any concentration in the entire temperature range at which the liquid mixture exists. The refractive index of g-picoline at a wavelength of  $\lambda = 589.3$  nm is 1.5, so it should be expected that in the concentration range for which an anomalous light scattering peak is recorded, a nanodroplet phase enriched in g-picoline can be observed, and the phase contrast, in this case, will be high enough to study these nanodroplets with a phase microscope, see [7 - 10]; we are currently experimenting with. In addition, in [12 - 15] it was shown that the Landau-Placek relation

$$\left( \frac{I_{\omega}}{I_{\omega-\delta\omega} + I_{\omega+\delta\omega}} \right)$$
, where  $I_{\omega}$  - unbiased component intensity,  $I_{\omega\pm d\omega}$  - the intensity of the displaced components increases with decreasing solution concentration over a wide temperature range. In [16], it was suggested that the increase in 
$$\left( \frac{I_{\omega}}{I_{\omega-\delta\omega} + I_{\omega+\delta\omega}} \right)$$
 is due to an increase in the isothermal compressibility  $b_T$  with an increase in the water content. This, in turn, may be due to the appearance of stable supramolecular structures (nuclei of the nanodroplet phase) with a decrease in the concentration of g-picoline; this phase, in accordance with the model presented in [7 - 10], causes anomalous scattering in an aqueous solution of g-

picoline at its concentration of 0.06 ppm. in solution; henceforth, we will call this concentration the critical concentration. This paper presents the results of a study of the frequency dispersion of hypersound velocity in aqueous solutions of g-picoline in a wide range of temperatures and concentrations; special attention is paid to the concentration region near the critical one.

### Experiment.

The experimental setup is described in detail in [17]. The spectra of polarized light scattering were studied by using a two-pass Fabry-Perot interferometer. The light scattering angle was  $\sim 90^\circ$  and  $135^\circ$ . The error in setting the angle did not exceed  $0.2^\circ$ . The dispersion region of the interferometer was  $60.625\text{cm}^{-1}$  and  $60.417\text{cm}^{-1}$ . The interference contrast of the pattern reached  $4 \cdot 10^5$ , the sharpness of the interference maxima was 35. The radiation source was a He-Ne laser with a wavelength of 632.8 nm and a power of 115 mW. The shift of the Mandelstam - Brillouin components  $D_{\omega}$  was measured with an error not exceeding 0.5%. The hypersound velocity  $v$  was calculated using the formula:

$$v = \frac{\Delta\omega \cdot c}{\omega_0 2n \sin \frac{\theta}{2}}, \quad (1)$$

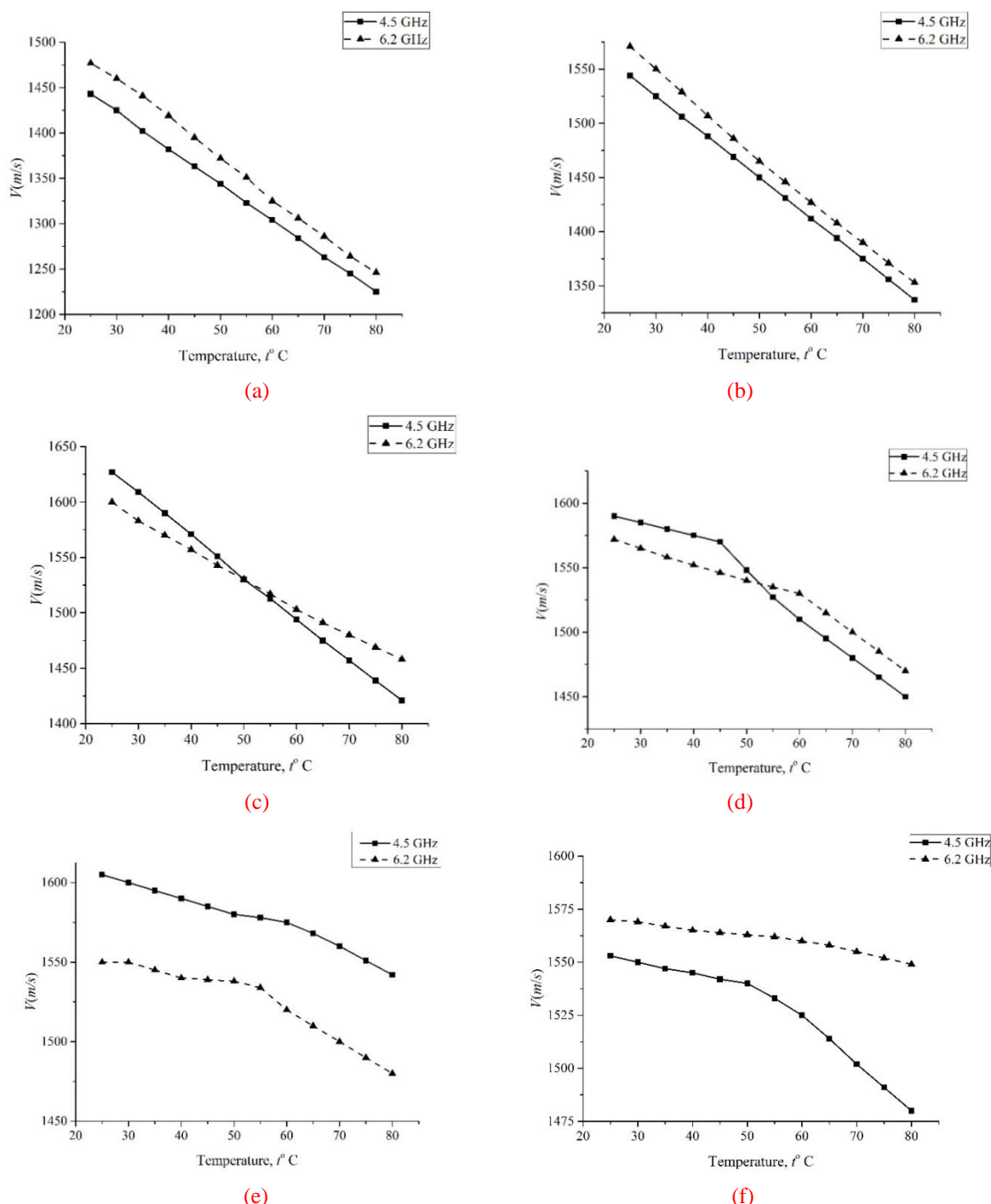
where  $c$  - is the speed of light,  $\omega_0$  - is the frequency of the exciting light,  $n$  - is the refractive index,  $\theta$  - is the scattering angle.

### Experimental results and discussion.

On Fig. 1 (a) - (f) panels show the temperature dependences of the hypersound velocity  $v$  for frequencies of 6.2 GHz and 4.5 GHz for pure g-picoline (panel (a)) and for concentrations of g-picoline  $C = 0.4$  ppm (panel (b));  $C = 0.2$  ppm (panel (c));  $C = 0.1$  ppm (panel (d));  $C = 0.06$  ppm (panel (e)) and  $C = 0.05$  ppm (panel (f)). Figure 2 shows the dependence of the area difference  $D_S$  under the curves for frequencies of 6.2 GHz and 4.5 GHz on the concentration  $C$  of the solution.

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**Fig. 1.** Dependences of the velocity  $V$  of hypersound at frequencies 4.5 and 6.2. GHz on temperature for pure g-picoline (a), mixture "g picoline - water" at a concentration of g-picoline 0.4 mole fractions (b), mixture "g-picoline - water" at a concentration of g- picoline 0.2 mole fractions (c), mixtures "g-picoline - water" at a concentration of g-picoline 0.1 mole fractions (d), mixtures "g- picoline - water" at a concentration of g- picoline 0.06 mole fractions (e), mixtures "g-picoline - water" at a concentration of g- picoline 0.05 mole fraction (f).

As can be seen from the figure (panels (a) and (b)), the nature of the temperature dependence at the indicated frequencies for pure g-picoline and 0.4 ppm is approximately the same for both frequencies. Differences appear when the concentration of g-

picoline is  $C = 0.2$  ppm and 0.1 ppm (panels (c) and (d)). It can be seen that at  $t > 50$  °C, there is a positive and at  $t < 50$  °C, a negative hypersonic speed dispersion in the frequency interval between 6.2 and 4.5 GHz.

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Figure 2 shows the dependence of the area difference  $DS = S_{6.2} - S_{4.5}$ , where  $S_{6.2}$  and  $S_{4.5}$  are the areas under the graphs in Fig. 1 for frequencies of 6.2 and 4.5 GHz, respectively, depending on the concentration of g-picoline. Since  $DS$  is the

temperature-averaged difference in hypersonic speeds at the specified frequencies, this value can be conditionally considered as a quantitative characteristic of the hypersound velocity dispersion.

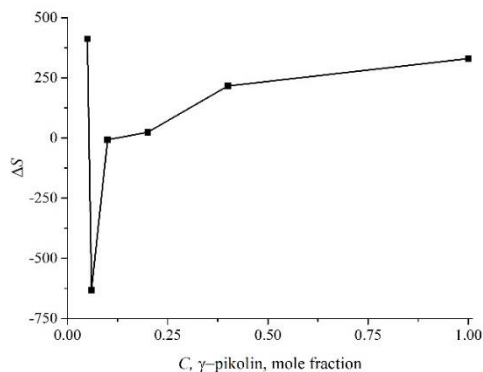


Fig. 2. Concentration dependence of  $DS$

In this case, the concentration dependence of  $DS$  anticorrelates with the concentration dependence of the scattering coefficient (see [1]): a sharp negative minimum of  $DS$  at  $C = 0.06$  ppm. coincides with the position of the scattering peak at the same concentration. Thus, it has been found that the appearance of a nanodroplet phase (peak of anomalous scattering) leads to a change in the sign of the dependence  $dV/dw$ , where  $V$  - is the hypersound velocity.

### Conclusions

In aqueous solutions of g-picoline, the situation occurs when the bulk viscosity relaxes at a relatively low frequency compared to other aqueous solutions. This led to the fact that at frequencies  $f > 4$  GHz, hypersonic speed dispersion appeared, which was not associated with the processes of relaxation of bulk viscosity. The temperature, concentration, and frequency dependences of the hypersound velocity consistently reflect the following: a) Positive sound

velocity dispersion for pure g- picoline and for a solution with a concentration of  $C=0.4$  ppm. g-p. b) Positive dispersion for  $C=0.2$  and  $0.1$  ppm. g-p. for temperature  $t > 50^\circ\text{C}$ , negative dispersion for temperatures  $t < 50^\circ\text{C}$ , c) For a concentration of  $0.06$  ppm g-p. only negative variance is observed. Such a transformation of the sign of the dispersion can be associated with the process of solution structurization, which begins at a concentration of  $0.2$  ppm. g-p. and reaches its maximum value at  $C=0.06$  ppm. g-p. In parallel with structurization, the thermodynamic state of the solution approaches (ECP) or (LCP) [16].

With a decrease in the concentration of g-picoline, the solution is closer to the periphery of the closed separation curve, and therefore the influence of the Upper critical point and the Lower critical point becomes weaker. With a decrease in the concentration of g-picoline in water, the process of structurization intensifies and its influence increases to a maximum associated with the optimal number of g-picoline molecules.

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